

*Defects, dualities and gauging in string theory
via gerbes*

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(based, in part, on joint work with Gawędzki & Waldorf, with Runkel)

Nordic String Theory Meeting 2010

Leibniz Universität Hannover, 22–23 February 2010

Part I

A 2-category for the 2D σ -model

The gerbe for the monophasic world-sheet

Lagrangian description of the (critical) string – the simplest scenario

$$X \quad : \quad \begin{array}{ccc} \begin{array}{c} \text{[Image of a green torus with three holes]} \\ (\Sigma, \gamma) \end{array} & \mapsto & \begin{array}{c} \text{[Image of a blue mesh surface with three holes]} \\ (M, g, \mathcal{G}), \text{ curv}(\mathcal{G}) =: H \in Z^3(M) \end{array} \end{array}$$

governed by the action functional $(dX = \underset{\text{loc.}}{\partial_a X^\mu} d\sigma^a \otimes \partial_\mu)$

$$S_\sigma[X; \gamma] = -\frac{1}{2} \int_\Sigma g(dX \wedge \star_\gamma dX) - i \log \text{Hol}_\mathcal{G}(X)$$

$\text{Hol}_\mathcal{G}(X)$ is the **SURFACE HOLONOMY**

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$\text{Hol}_\mathcal{G}(X)$ is the **SURFACE HOLONOMY** of **GERBE** \mathcal{G} , locally given by

$$\left\{ \begin{array}{l} H|_{\mathcal{O}_i} =: dB_i \\ (B_j - B_i)|_{\mathcal{O}_i \cap \mathcal{O}_j} =: dA_{ij} \\ (A_{jk} - A_{ik} + A_{ij})|_{\mathcal{O}_i \cap \mathcal{O}_j \cap \mathcal{O}_k} =: i d \log g_{ijk} \\ (g_{jkl} \cdot g_{ikl}^{-1} \cdot g_{ijl} \cdot g_{ijk}^{-1})|_{\mathcal{O}_i \cap \mathcal{O}_j \cap \mathcal{O}_k \cap \mathcal{O}_l} = 1 \end{array} \right. \quad \text{mod} \quad \left\{ \begin{array}{l} B_i \mapsto B_i + d\Pi_i \\ A_{ij} \mapsto A_{ij} + (\Pi_j - \Pi_i)|_{\mathcal{O}_i \cap \mathcal{O}_j} - i d \log \chi_{ij} \\ g_{ijk} \mapsto g_{ijk} \cdot (\chi_{jk}^{-1} \cdot \chi_{ik} \cdot \chi_{ij}^{-1})|_{\mathcal{O}_i \cap \mathcal{O}_j \cap \mathcal{O}_k} \end{array} \right.$$

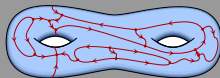
Pour qu'on n'en ait pas (que) la gerbe...



La gerbe, Henri Matisse (1953)

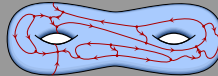
The string background for the multiphase world-sheet

Natural generalisation (e.g., strings on orbifolds and T-folds):
 Σ with embedded **DEFECT** Γ

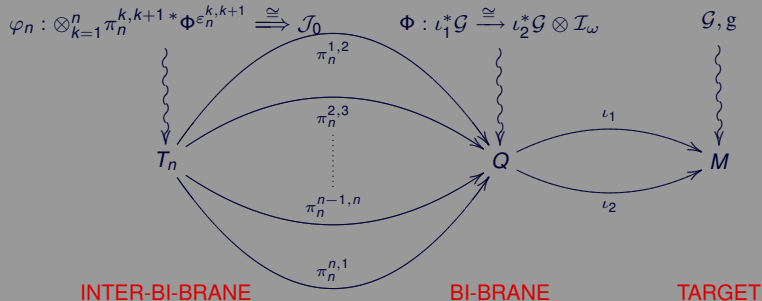


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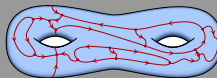


σ -model requires **STRING BACKGROUND** $\mathfrak{B} = (\mathcal{M}, \mathcal{B}, \mathcal{J})$

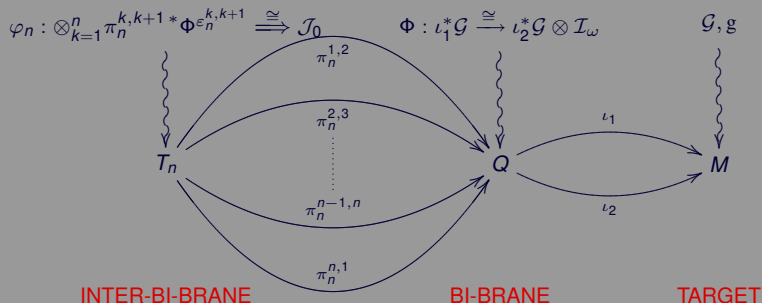


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Example: D-brane \equiv boundary bi-brane for $M = \mathcal{M} \sqcup \{\bullet\}$,

$$\iota_1 : Q = \mathcal{D} \hookrightarrow \mathcal{M}, \quad \iota_2 : \mathcal{D} \rightarrow \{\bullet\}, \quad \Phi_{\mathcal{D}} : \mathcal{G}|_{\mathcal{D}} \xrightarrow{\cong} \mathcal{I}_{\omega=B_i+dA_i}$$

The string background for the multiphase world-sheet – ctd.

Upshot: cohomological classification scheme for σ -models

FIELD THEORY	CFT_α	$\text{D}_{\alpha\beta}$	$\text{J}_{\alpha_1\alpha_2\dots\alpha_n}$
GEOMETRY	TARGET $\mathcal{M} = (M, g, \mathcal{G})$	BI-BRANE $\mathcal{B} = (Q, \omega, \iota_1, \iota_2, \Phi)$	INTER-BI-BRANE $\mathcal{I} = (T_n, \varphi_n, (\varepsilon_n^{k,k+1}, \pi_n^{k,k+1}) n \in \mathbb{N}_{>0})$
2-CATEGORY	object	1-morphism	2-morphism
$\mathfrak{B}\mathfrak{C}\mathfrak{a}\mathfrak{t}(\mathcal{M} \sqcup \mathcal{Q} \sqcup \mathcal{T})$	\mathcal{G}	$\Phi : \iota_1^* \mathcal{G} \xrightarrow{\sim} \iota_2^* \mathcal{G} \star \mathcal{I}(\omega)$	$\varphi_n : \circ_{k=1}^n \pi_n^{k,k+1*} \Phi^{\varepsilon_n^{k,k+1}} \xrightarrow{\sim} \text{id}$
classification	$W^2(M, \mathbb{H})$ is $H^2(M, \text{U}(1))$ -torsor	$W^1(M, \mathcal{G}_1 \rightarrow \mathcal{G}_2)$ is $H^1(M, \text{U}(1))$ -torsor	$W^0(M, \Phi_1 \Rightarrow \Phi_2)$ is $\text{U}(1)^{ \pi_0(M) }$ -torsor

Part II

Dualities via world-sheet defects

The canonical interpretation of defects – lines

Categorical quantisation & more geometric analyses suggest

(some) DEFECTS \sim STRING DUALITIES

This can be rendered rigorous in the 2-categorical setting. . .

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Thm.: \mathfrak{B} canonically defines PREQUANTUM BUNDLE

$$\mathcal{L}_\sigma \rightarrow \mathbf{P}_\sigma, \quad \text{curv}(\mathcal{L}_\sigma) = \Omega_\sigma$$

Def.: DUALITY $\equiv \Omega_\sigma^-$ -lagrangean submanifold

$$\mathfrak{D}_\sigma \subset \mathbf{P}_\sigma \times \mathbf{P}_\sigma, \quad \text{pr}_1^* \mathcal{H}_\sigma = \text{pr}_2^* \mathcal{H}_\sigma$$

together with a bundle isomorphism

$$\text{pr}_1^* \mathcal{L}_\sigma|_{\mathfrak{D}_\sigma} \xrightarrow{\cong} \text{pr}_2^* \mathcal{L}_\sigma|_{\mathfrak{D}_\sigma}$$

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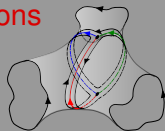
$$\text{pr}_1^* \mathcal{L}_\sigma|_{\mathfrak{D}_\sigma} \xrightarrow{\cong} \text{pr}_2^* \mathcal{L}_\sigma|_{\mathfrak{D}_\sigma}$$

Thm.: \mathcal{B} canonically defines a duality iff

- $\tilde{\iota}_\alpha : \text{LQ} \rightarrow \text{LM} : X \mapsto \iota_\alpha \circ X$ are **surjective submersions**
- (Γ, X) **topological**
- extra conditions (technical)

The canonical interpretation of defects – junctions

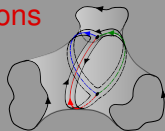
Defect junctions naturally associated with **interactions**



Further hints from categorial quantisation & study of **maximally symmetric WZW defects**.

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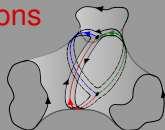
$$\mathfrak{I}_\sigma(\otimes \mathcal{B} : \mathcal{J} : \mathcal{B}) \subset \mathbb{P}_\sigma \times \mathbb{P}_\sigma \times \mathbb{P}_\sigma$$

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Example: 2-iso's for maxym G-WZW defects \sim

spaces of conformal blocks on punctured and decorated $\mathbb{C}P^1$ via $\mathrm{CS}_k(G)$ on $\mathbb{R} \times \mathbb{C}P^1_{\{P_k\}_{k \in \overline{1, n}}}$ with parallel Wilson lines of fixed holonomy

Symmetries as distinguished dualities

Converse result for a class of dualities with local data

$$\Phi_\sigma : \text{pr}_1^* \mathcal{L}_\sigma |_{\mathcal{D}_\sigma} \xrightarrow{\cong} \text{pr}_2^* \mathcal{L}_\sigma |_{\mathcal{D}_\sigma},$$

$$i \log \Phi_{\sigma i}[(X_1, p_1), (X_2, p_2)] = \int_{S^1} \text{Vol}(S^1) p_{2\mu} F^\mu(X_1) + W_i[(X_1, X_2)]$$

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Thm.: Φ_σ can. defines a flat bi-brane with

- world-volume $Q = (\text{id}_M \times F)(M) \subset M \times M$, with $F \in \text{Isom}(M, g)$;
- bi-brane maps $\iota_\alpha = \text{pr}_\alpha$, $\alpha \in \{1, 2\}$;
- bi-brane 1-isomorphism $\Phi : \mathcal{G} \xrightarrow{\cong} F^* \mathcal{G}$

These are **INTERNAL SYMMETRIES** of the closed string.

Part III

Generalised geometry with a 2-categorical twist

Brackets on the state space and on the target space

Observation: GENERALISED GEOMETRY natural in the symplectic setting (P, Ω) , via HAMILTONIAN SECTIONS:

$$\mathfrak{X}_h = \mathcal{X}_h \oplus h \in \ker d\Omega \subset \Gamma(E^{(1,0)}P), \quad E^{(1,0)}P := \wedge^1 TP \oplus \wedge^0 T^*P \rightarrow P$$

and Ω -TWISTED VINOGRADOV BRACKET:

$$[\mathfrak{X}_{h_1}, \mathfrak{X}_{h_2}]_{\mathbb{V}}^{\Omega} := [\mathcal{X}_{h_1}, \mathcal{X}_{h_2}] \oplus (\mathcal{X}_{h_1} \lrcorner dh_2 - \mathcal{X}_{h_2} \lrcorner dh_1 + \mathcal{X}_{h_1} \lrcorner \mathcal{X}_{h_2} \lrcorner \Omega) \equiv \mathfrak{X}_{\{h_1, h_2\}_{\Omega}}$$

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Idea: Given the manifold structure on the space of σ -model fields, we can look for bracket structures on $E^{(1,\bullet)}(M \sqcup Q \sqcup T) \rightarrow M \sqcup Q \sqcup T$ closing on σ -SYMMETRIC SECTIONS, with a homomorphic lift to

$$\text{CANONICAL} \\ \text{VINOGRADOV STRUCTURE} \quad \mathfrak{V}^\Omega P_\sigma := (E^{(1,0)}P_\sigma, [\cdot, \cdot]_V^{\Omega_\sigma}, \alpha_{TP_\sigma})$$

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Hint: The answer is known for $P_{\sigma, \emptyset}$: Courant algebroid $\mathfrak{C}^H M$ on $E^{(1,1)}M$ with Courant bracket twisted by H à la Ševera–Weinstein, Hitchin-isomorphic with $\mathfrak{C}_{\mathcal{G}} M$

Brackets on the state space and on the target space – ctd.

Thm.: In the presence of defects, the answer given by $(\Delta_Q := \iota_2^* - \iota_1^*)$

$$\mathfrak{M}^{(1,0),(H,\omega;\Delta_Q)}(M \sqcup Q) := (E^{(1,1)}M \sqcup E^{(1,0)}Q, [[\cdot, \cdot]]^{(H,\omega;\Delta_Q)}, (\cdot, \cdot)_{\sqcup}, \alpha_{T(M \sqcup Q)})$$

with **TWISTED BRACKET** on $\mathfrak{Y}_i = ({}^M\mathfrak{Y}_i, {}^Q\mathfrak{Y}_i) = ({}^M\mathcal{Y}_i \oplus v_i, {}^Q\mathcal{Y}_i \oplus \xi_i)$

$$[[\mathfrak{Y}_1, \mathfrak{Y}_2]]^{(H,\omega;\Delta_Q)}|_M = [{}^M\mathcal{Y}_1, {}^M\mathcal{Y}_2] \oplus (\mathcal{L}_{{}^M\mathcal{Y}_1} v_2 - \mathcal{L}_{{}^M\mathcal{Y}_2} v_1 - \frac{1}{2} d({}^M\mathcal{Y}_1 \lrcorner v_2 - {}^M\mathcal{Y}_2 \lrcorner v_1) + {}^M\mathcal{Y}_1 \lrcorner {}^M\mathcal{Y}_2 \lrcorner H),$$

$$[[\mathfrak{Y}_1, \mathfrak{Y}_2]]^{(H,\omega;\Delta_Q)}|_Q = [{}^Q\mathcal{Y}_1, {}^Q\mathcal{Y}_2] \oplus ({}^Q\mathcal{Y}_1 \lrcorner d\xi_2 - {}^Q\mathcal{Y}_2 \lrcorner d\xi_1 + {}^Q\mathcal{Y}_1 \lrcorner {}^Q\mathcal{Y}_2 \lrcorner \omega + \frac{1}{2} ({}^Q\mathcal{Y}_1 \lrcorner \Delta_Q v_2 - {}^Q\mathcal{Y}_2 \lrcorner \Delta_Q v_1)),$$

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Observation: Study of **automorphisms** and **Hitchin-type iso's**

$$\mathfrak{M}_{\iota_\alpha}^{(1,0),(H,\omega;\Delta_Q)}(M \sqcup Q) \cong \mathfrak{M}_{(\mathcal{G}, \mathcal{B}), \iota_\alpha}^{(1,0),(0,0;\Delta_Q)}(M \sqcup Q)$$

indicate that generalised geometry is a natural generalisation of the geometry of TM in the presence of $\mathfrak{B}\mathfrak{G}\mathfrak{r}\mathfrak{b}^\nabla(M \sqcup Q \sqcup T)$

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In keeping with the above,

Thm.: \mathcal{B} can. induces a **morphism** in the category of twisted Courant algebroids on $E^{(1,1)}Q$.

Gaugeability constraints for σ -model symmetries

Internal symmetries of $S_\sigma \equiv \sigma$ -SYMMETRIC ι_α -ALIGNED SECTIONS

$$\mathfrak{K}_a = ({}^M\mathfrak{K}_a, {}^Q\mathfrak{K}_a) = ({}^M\mathcal{K}_a \oplus \kappa_a, {}^Q\mathcal{K}_a \oplus k_a), \quad \left\{ \begin{array}{l} \iota_\alpha^* {}^Q\mathcal{K}_a = {}^M\mathcal{K}_a|_{\iota_\alpha(Q)} \\ d_H {}^M\mathfrak{K}_a = 0 \\ d_\omega {}^Q\mathfrak{K}_a + \Delta_Q \kappa_a = 0 \end{array} \right.$$

$$a \in \overline{1, \dim \mathfrak{k}_{\sigma, \iota_\alpha}}$$

The corresponding **hamiltonian sections**

$$\tilde{\mathfrak{K}}_a = e^{\text{pr}_{T^*L_\sigma M} \theta_{T^*L_\sigma M}} \triangleright \tilde{L}\mathfrak{K}_a, \quad \tilde{L} : \Gamma_{\iota_\alpha}(E^{(1,1)}M \sqcup E^{(1,0)}Q) \rightarrow \Gamma(E^{(1,0)}P_\sigma),$$

written in terms of the canonical 1-form $\theta_{T^*L_\sigma M} \in \Omega^1(T^*L_\sigma M)$, obey

$$\left[\tilde{\mathfrak{K}}_a, \tilde{\mathfrak{K}}_b \right]_{\mathbb{V}}^{\Omega_\sigma} = \left[[\mathfrak{K}_a, \widetilde{\mathfrak{K}}_b] \right]^{(H, \omega; \Delta_Q)}$$

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The corresponding **hamiltonian sections**

$$\tilde{\mathfrak{K}}_a = e^{\text{pr}_{T^*L_\sigma M} \theta_{T^*L_\sigma M}} \triangleright \tilde{L}\mathfrak{K}_a, \quad \tilde{L} : \Gamma_{\iota_\alpha}(E^{(1,1)}M \sqcup E^{(1,0)}Q) \rightarrow \Gamma(E^{(1,0)}P_\sigma),$$

written in terms of the canonical 1-form $\theta_{T^*L_\sigma M} \in \Omega^1(T^*L_\sigma M)$, obey

$$\left[\tilde{\mathfrak{K}}_a, \tilde{\mathfrak{K}}_b \right]_{\mathbb{V}}^{\Omega_\sigma} = \left[\mathfrak{K}_a, \widetilde{\mathfrak{K}}_b \right]^{(H, \omega; \Delta_Q)}$$

The realisation of $\mathfrak{k}_{\sigma, \iota_\alpha}$ on the state space becomes **hamiltonian** iff

$$\left[\mathfrak{K}_a, \mathfrak{K}_b \right]^{(H, \omega; \Delta_Q)} = f_{ab}^c \mathfrak{K}_c$$

Gaugeability constraints for σ -model symmetries – ctd.

Conclusion: Necessary conditions of gaugeability of $\mathfrak{k}_{\sigma, \iota_\alpha}$:

$$\left(\bigoplus_{a \in \overline{1, \dim \mathfrak{k}_{\sigma, \iota_\alpha}}} \mathbb{R} \mathfrak{K}_a, [[\cdot, \cdot]]^{(H, \omega; \Delta_Q)} \right) \cong \mathfrak{k}_{\sigma, \iota_\alpha} \quad \wedge \quad (\mathfrak{K}_a, \mathfrak{K}_b)_\perp = 0$$

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N.B. The ‘isotropy’ condition nullifies the **anomaly** of the Poisson algebra of the Noether currents.

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Conclusion: Necessary conditions of gaugeability of $\mathfrak{k}_{\sigma, \ell_\alpha}$:

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N.B. The ‘isotropy’ condition nullifies the **anomaly** of the Poisson algebra of the Noether currents.

Equivalently, the gaugeability relations ensure the existence of a **$\mathfrak{k}_{\sigma, \ell_\alpha}$ -equivariantly closed extension** \widehat{H} of H , and a **$\mathfrak{k}_{\sigma, \ell_\alpha}$ -equivariant extension** $\widehat{\omega}$ of ω in the Cartan model of $\mathfrak{k}_{\sigma, \ell_\alpha}$ -equivariant cohomology of $M \sqcup Q$, s.t.

$$\widehat{d}\widehat{H} = 0, \quad \widehat{d}\widehat{\omega} = -\Delta_Q \widehat{H}$$

for $\widehat{d}\eta(X) = d\eta(X) + X^a \mathcal{K}_{a\lrcorner} \eta(X)$, $X \in \mathfrak{k}_{\sigma, \ell_\alpha}$.

Part IV

The gauged σ -model

Motivation for & problems with gauging

I Motivation:

- I.1 string theory on cosets (via gauging & symplectic reduction);
- I.2 T-duality etc.

II Problems:

- II.1 lifting the geometric action of the isometry (sub-)group to $\mathcal{B}\mathcal{G}\text{rb}^\nabla(M \sqcup Q \sqcup T)$;
- II.2 coupling (non-trivial) world-sheet gauge fields to $\mathcal{B}\mathcal{G}\text{rb}^\nabla(M \sqcup Q \sqcup T)$.

Equivariant structures

Simplicial descent schemes: in the absence of defects,

$$\begin{array}{ccccccc}
 \cdots & \xrightarrow{\text{Gd}_i^{(4)}} & \text{G}^3 \times M & \xrightarrow{\text{Gd}_i^{(3)}} & \text{G}^2 \times M & \xrightarrow{\text{Gd}_i^{(2)}} & \text{G} \times M & \xrightarrow{\text{Gd}_i^{(1)}} & M \left(\begin{array}{c} - \\ - \\ \xrightarrow{M \wr G} M/G \end{array} \right) \\
 & & \delta_G \gamma = 1 & & \gamma & & \Upsilon & & \mathcal{G} & & \bar{\mathcal{G}} \\
 & & \text{associativity} & & \text{distributiveness} & & \text{element-wise} & & & &
 \end{array}$$

where

$$\Upsilon : \text{Gd}_1^{(1)*} \mathcal{G} \xrightarrow{\cong} d_0^{(1)*} \mathcal{G} \otimes \mathcal{I}_\rho, \quad \gamma : (\text{Gd}_0^{(2)*} \Upsilon \otimes \text{id}) \circ \text{Gd}_2^{(2)*} \Upsilon \xrightarrow{\cong} \text{Gd}_1^{(2)*} \Upsilon$$

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Observation: Gaugeability **fixes** $\rho = \theta_L^a \wedge \kappa_a + \frac{1}{2} \theta_L^a \wedge \theta_L^b ({}^M \mathcal{K}_{a \lrcorner} \kappa_b)$

Equivariant structures – ctd.

Similarly, in the presence of defects,

$$\begin{array}{ccccccc}
 \cdots & \xrightarrow{\text{G}d_i^{(3)}} & G^2 \times Q & \xrightarrow{\text{G}d_i^{(2)}} & G \times Q & \xrightarrow{\text{G}d_i^{(1)}} & Q \left(- \xrightarrow{\alpha_{\mathcal{W}_G}} Q/G \right), \\
 & & \text{c.c.}(\Xi, \gamma) & & \Xi & & \Phi & & \bar{\Phi}
 \end{array}$$

where

$$\Xi : (\iota_2^* \Upsilon \otimes \text{id}) \circ \circ d_1^{(1)*} \Phi \xrightarrow{\cong} (\circ d_0^{(1)*} \Phi \otimes \mathcal{J}_\lambda) \circ \iota_1^* \Upsilon$$

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Equivariant structures – ctd.

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Observation: Gaugeability **fixes** $\lambda = \theta_L^a k_a$

‘Problem’:

$$\text{Thm.:} \quad \left\{ \begin{array}{l} \text{Equivalence classes} \\ \text{of } (G, 0)\text{-equivariant bi-branes} \\ \text{with world-volume } Q \\ \text{for } (G, 0)\text{-equivariant gerbes over } M \end{array} \right\} \xleftrightarrow{1:1} \left\{ \begin{array}{l} \text{Equivalence classes} \\ \text{of bi-branes} \\ \text{with world-volume } Q/G \\ \text{for gerbes over } M/G \end{array} \right\}.$$

However, $k_a \neq 0$ generically.

The coupling of the world-sheet gauge field

Observation: Gauging G prerequisites replacing $X \in C^1(\Sigma, M \sqcup Q)$ with $X \in \Gamma(P \times_G (M \sqcup Q))$ for

principal G -bundle $G \hookrightarrow P \rightarrow \Sigma$, with principal G -connection $\mathcal{A} \in \Omega^1(P) \otimes \mathfrak{g}$

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The study of the conditions of invariance (for P trivial) of

$$S_{\sigma, \text{top}}[(\Gamma, X)] \mapsto S_{\sigma, \text{top}}[(\Gamma, X)] + \int_{\Sigma} (X \times \text{id}_{\Sigma})^* \zeta_{\mathcal{A}} + \int_{\Gamma} ((X \times \text{id}_{\Gamma})|_{\Gamma})^* \mu_{\mathcal{A}},$$

$$\zeta_{\mathcal{A}}(\sigma, m) = -\alpha_a(m) \wedge A^a(\sigma) + \frac{1}{2} \beta_{ab}(m) A^a \wedge A^b(\sigma), \quad \mu_{\mathcal{A}}(\sigma, m) = \gamma_a(m) A^a(\sigma)$$

leads to the definitions

$$\mathcal{G}_{\mathcal{A}} := \text{pr}_2^* \mathcal{G} \otimes \mathcal{I}_{\rho_{\mathcal{A}}}, \quad \rho_{\mathcal{A}}(\rho, m) := -\kappa_a(m) \wedge \mathcal{A}^a(\rho) + \frac{1}{2} ({}^M \mathcal{X}_a \lrcorner \kappa_b)(m) \mathcal{A}^a \wedge \mathcal{A}^b(\rho),$$

$$\Phi_{\mathcal{A}} := \text{pr}_2^* \Phi \otimes \mathcal{J}_{\lambda_{\mathcal{A}}}, \quad \lambda_{\mathcal{A}}(\rho, m) := k_a(m) \mathcal{A}^a(\rho)$$

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$$\Phi_{\mathcal{A}} := \text{pr}_2^* \Phi \otimes \mathcal{J}_{\lambda_{\mathcal{A}}}, \quad \lambda_{\mathcal{A}}(\rho, m) := k_a(m) \mathcal{A}^a(\rho)$$

Thm.: $\mathcal{G}_{\mathcal{A}}$ carries a **canonical structure** of a $(G, 0)$ -equivariant gerbe on $P \times M$, and $\Phi_{\mathcal{A}}$ carries a **canonical structure** of a $(G, 0)$ -equivariant $\mathcal{G}_{\mathcal{A}}$ -bi-brane on $P|_{\Gamma} \times Q$.

The coupling of the world-sheet gauge field – ctd.

Corollary: $(\mathcal{G}_{\mathcal{A}}, \Phi_{\mathcal{A}})$ descend to unique (equivalence classes of) $(\overline{\mathcal{G}}_{\mathcal{A}}, \overline{\Phi}_{\mathcal{A}})$ over $P \times_G (M \sqcup Q)$, and so can be used to define the **G-GAUGED σ -MODEL**

$$S_{\sigma}[(\Gamma, X); \mathcal{A}, \gamma] = S_{\sigma, \text{kin}}[X; \mathcal{A}, \gamma] - i \log \text{Hol}_{\overline{\mathcal{G}}_{\mathcal{A}}, \overline{\Phi}_{\mathcal{A}}}(\Gamma, X),$$

with $S_{\sigma, \text{kin}}[X; \mathcal{A}, \gamma]$ obtained through **minimal coupling**.

The coupling of the world-sheet gauge field – ctd.

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Recall that the **gauge group** G_{Σ} is the set $\Gamma(P \times_{\text{Ad}G} G)$ with the group operation induced from

$$[(p, g_1)] \cdot [(p, g_2)] := [(p, g_1 \cdot g_2)]$$

and with the action on P induced by

$$(P \times_{\text{Ad}G} G) \times P \rightarrow P : ([(\tau_i(\sigma, g), h)], \tau_i(\sigma, g)) \mapsto \tau_i(\sigma, h \cdot g) = : [(\tau_i(\sigma, g), h)] \triangleright \tau_i(\sigma, g).$$

as per

$$\lambda. : \Gamma(P \times_{\text{Ad}G} G) \times P \rightarrow P : (\chi_i, \tau_i(\sigma, g)) \mapsto \chi_i(\sigma) \triangleright \tau_i(\sigma, g) = : \lambda_{(\chi_i)}(\tau_i(\sigma, g)).$$

The coupling of the world-sheet gauge field – ctd.

We have the fundamental

Thm.: $S_\sigma[(\Gamma, X); \mathcal{A}, \gamma]$ is **invariant** under G-gauge transformations

$$((\chi_i), X) \mapsto ((\lambda_{(\chi_i)} \circ \text{pr}_1) \times \text{pr}_2) \circ X, \quad \mathcal{A} \mapsto \lambda_{(\chi_i^{-1})}^* \mathcal{A}.$$

Proof: uses the G-equivariance of \mathcal{A} and the form of $\rho_{\mathcal{A}}$ and $\lambda_{\mathcal{A}}$.

Part V

Outlook

Outlook

- understanding T-duality, with particular emphasis on geometric structures behind the metric, the torsion and the dilaton;
- construction of spaces modelled on toroidal bundles only locally;
- including supersymmetry in the generalised geometric framework with a 2-categorical twist;
- study of the effective gerbe-twisted gauge field theory and the emergent geometry of bi-branes in the gerbe-theoretic context;
- gerbe theory vs criticality (generalised Ricci flows?);
- ‘holographic principle’ for higher categorial structures;
- ...